

# Analytic Solutions for Pricing Double Barrier Options in the Presence of Stochastic Volatility

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# Task

- There exist closed-form solutions for vanilla options in the presence of stochastic volatility for nearly a decade (Heston, 1993)
- For double barrier options we still depend on slow numerical methods, in particular Finite Difference or Monte Carlo methods
- Lipton (2001) obtains semi-analytical solutions for a special case

# Overview

1. The Heston Model
2. Free Green's Function
3. Method of Images
4. Eigenfunction Expansion
5. Application to other options
6. Implementation

# 1. The Heston Model

## 1.1 Overview

- Variance of the spot process is no longer deterministic, but modelled as a second stochastic process
- Heston chose a mean-reverting square-root process, i.e.

$$dS_t = S_t(\mu dt + \sqrt{v(t)}dW_t^{(1)})$$
$$dv_t = \kappa(\theta - v(t))dt + \sigma\sqrt{v(t)}dW_t^{(2)}$$

where  $W_t^{(1)}$ ,  $W_t^{(2)}$  are Brownian Motions with correlation  $\rho$ .

## 1.2 Differential Equation

Every contingent claim  $U(t, v, S)$  paying  $g(S) = U(T, v, S)$  has to satisfy the partial differential equation

$$U_t + (r_d - r_f)SU_S + \frac{1}{2}\sigma^2vU_{vv} + \frac{1}{2}vS^2U_{SS} + \rho\sigma vSU_{vS} - r_dU + (\tilde{\kappa}(\tilde{\theta} - v) - \lambda v)U_v = 0$$

(Can be shown by creating a replication portfolio, if we assume the market price of volatility risk  $\lambda$  has the form  $\lambda(t, v, S) = \lambda v$ ).

After reverting the time, i.e. applying the transformation

$$\tau = T - t$$

and elimination of the market price of volatility risk  $\lambda$  by replacing  $\tilde{\kappa}$  and  $\tilde{\theta}$  by  $\kappa := \tilde{\kappa} + \lambda$  and  $\theta := \frac{\tilde{\kappa}\tilde{\theta}}{\tilde{\kappa} + \lambda}$  one can rewrite it as

$$U_\tau - (r_d - r_f)SU_S - \frac{1}{2}\sigma^2vU_{vv} - \frac{1}{2}vS^2U_{SS} - \rho\sigma vSU_{vS} + r_dU - \kappa(\theta - v)U_v = 0$$

## 1.3 Double Barrier Options

We concentrate on double-barrier knock-out calls, which have to satisfy the pde with conditions

$$\begin{aligned}U(T, v, S) &= (S - K)^+ \\U(t, v, A) &= U(t, v, B) = 0\end{aligned}$$

respectively

$$\begin{aligned}U(0, v, S) &= (S - K)^+ \\U(\tau, v, A) &= U(\tau, v, B) = 0\end{aligned}$$

## 2. Free Green's Function

## 2.1 Overview

- Transform the pricing problem into a more tractable form
- Then derive the transition probability density function for the problem without barriers.

## 2.2 Transformation

We want to write the derived pricing problem in forward terms to remove the interest rates.

Apply the transformation

$$\hat{U}(\tau, v, F) = \hat{U}(\tau, v, e^{(r_d - r_f)\tau} S) = U(\tau, v, S) e^{r_d \tau}$$

Results in

$$\hat{U}_\tau - \frac{1}{2}vF^2\hat{U}_{FF} - \rho\sigma vF\hat{U}_{vF} - \frac{1}{2}\sigma^2v\hat{U}_{vv} - \kappa(\theta - v)\hat{U}_v = 0$$

$$\hat{U}(0, v, F) = (F - K)^+$$

$$\hat{U}(\tau, v, \hat{A}) = \hat{U}(\tau, v, \hat{B}) = 0$$

$$\hat{U}_\tau(\tau, 0, F) - \kappa\theta\hat{U}_v(\tau, 0, F) = 0$$

with transformed barriers

$$\hat{A} = e^{(r_d - r_f)\tau} A, \quad \hat{B} = e^{(r_d - r_f)\tau} B$$

We now nondimensionalize the problem by representing  $\hat{U}$  as

$$\hat{U}(\tau, v, F) = K\Phi(\tau, v, \xi)$$

with  $\xi = F/K$  and then further reduce the problem via the transformation

$$\begin{aligned} \tau &\rightarrow \tau, & \xi &\rightarrow X = \ln(\xi) \\ \Phi(\tau, v, \xi) &\rightarrow W(\tau, v, X) = e^{-X/2}\Phi(\tau, v, \xi) \end{aligned}$$

This results in the following pricing problem:

$$W_\tau - \frac{1}{2}vW_{XX} - \frac{1}{2}\sigma^2vW_{vv} - \rho\sigma vW_{vX} - \hat{\kappa}(\hat{\theta} - v)W_v + \frac{1}{8}vW = 0$$

$$W(0, v, X) = (e^{X/2} - e^{-X/2}) +$$

$$W(\tau, v, \hat{a}) = W(\tau, v, \hat{b}) = 0$$

where

$$\hat{a} = \ln(\hat{A}/K), \quad \hat{b} = \ln(\hat{B}/K)$$

$$\hat{\kappa} = \kappa - \sigma\rho/2, \quad \hat{\theta} = \kappa\theta/\hat{\kappa}$$

One can show that this is the backward Kolmogoroff equation \* with killing at a rate of  $\frac{v}{g}$  for the following system of stochastic differential equations:

$$\begin{aligned}dX &= \sqrt{v}dW_t^{(1)} \\dv &= \hat{\kappa}(\hat{\theta} - v)dt + \sigma\sqrt{v}dW_t^{(2)}\end{aligned}$$

\*see Karatzas and Shreve (1991)

## 2.3 Derivation

Once we found the Green's function of our pricing problem, we can use the following theorem to obtain a solution:

**Theorem 1.** *Let  $p(t, X, v, t', X', v')$  be the Green's function of a process  $(X_t, v_t)$  and let  $w(X, v)$  be the payoff function. The pricing formula  $W(t, X, v)$  can then be expressed as*

$$W(t, v, X) = \int \int p(t, X, v, T, X', v') w(X', v') dX' dv'$$

In our case we can make use of the fact, that the coefficients and the killing rate are independent of  $t$ ,  $X$  and can hence concentrate on the differences  $\tau = t' - t$  and  $Y = X' - X$  rather than on the arguments themselves:

$$p(t, X, v, t', X', v') = p(\tau, Y, v, v')$$

Another simplification can be done, when we observe that the payoff of our double barrier call is independent of the volatility. Hence we only need to find its integral over  $v'$ , i.e.

$$q(\tau, Y, v) = \int_0^{\infty} p(\tau, Y, v, v') dv'$$

Due to the fact that the coefficients of  $q$  are linear functions of  $v$ , we can now guess the following form:

$$q(\tau, Y, v) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ikY + 2(\mathbf{A}(\tau, k) - \mathbf{B}(\tau, k)v)/\sigma^2} dk$$

After some calculations we obtain:

$$\mathbf{A}(\tau, k) = -\kappa\theta(\mu + \zeta)\tau - \kappa\theta \ln \left( \frac{-\mu + \zeta + (\mu + \zeta)e^{-2\zeta\tau}}{2\zeta} \right)$$
$$\mathbf{B}(\tau, k) = \frac{\sigma^2(k^2 + 1/4)(1 - e^{-2\zeta\tau})}{4(-\mu + \zeta + (\mu + \zeta)e^{-2\zeta\tau})}$$

where

$$\mu(k) = -\frac{1}{2}(ik\sigma\rho + \hat{\kappa})$$
$$\zeta(k) = \frac{1}{2}\sqrt{k^2\sigma^2(1 - \rho^2) + 2ik\sigma\rho\hat{\kappa} + \hat{\kappa}^2 + \frac{\sigma^2}{4}}$$

### 3. Method of Images

## 3.1 Overview

**Now:** Incorporate the barriers into the Free Green's function via the method of images.

**Principle:** Derive a generalized reflection principle and use it to express the restricted problem as an infinite sum of the unrestricted one.

**References:** Dewynne, Howison and Wilmott (1993); Andreasen (2001).

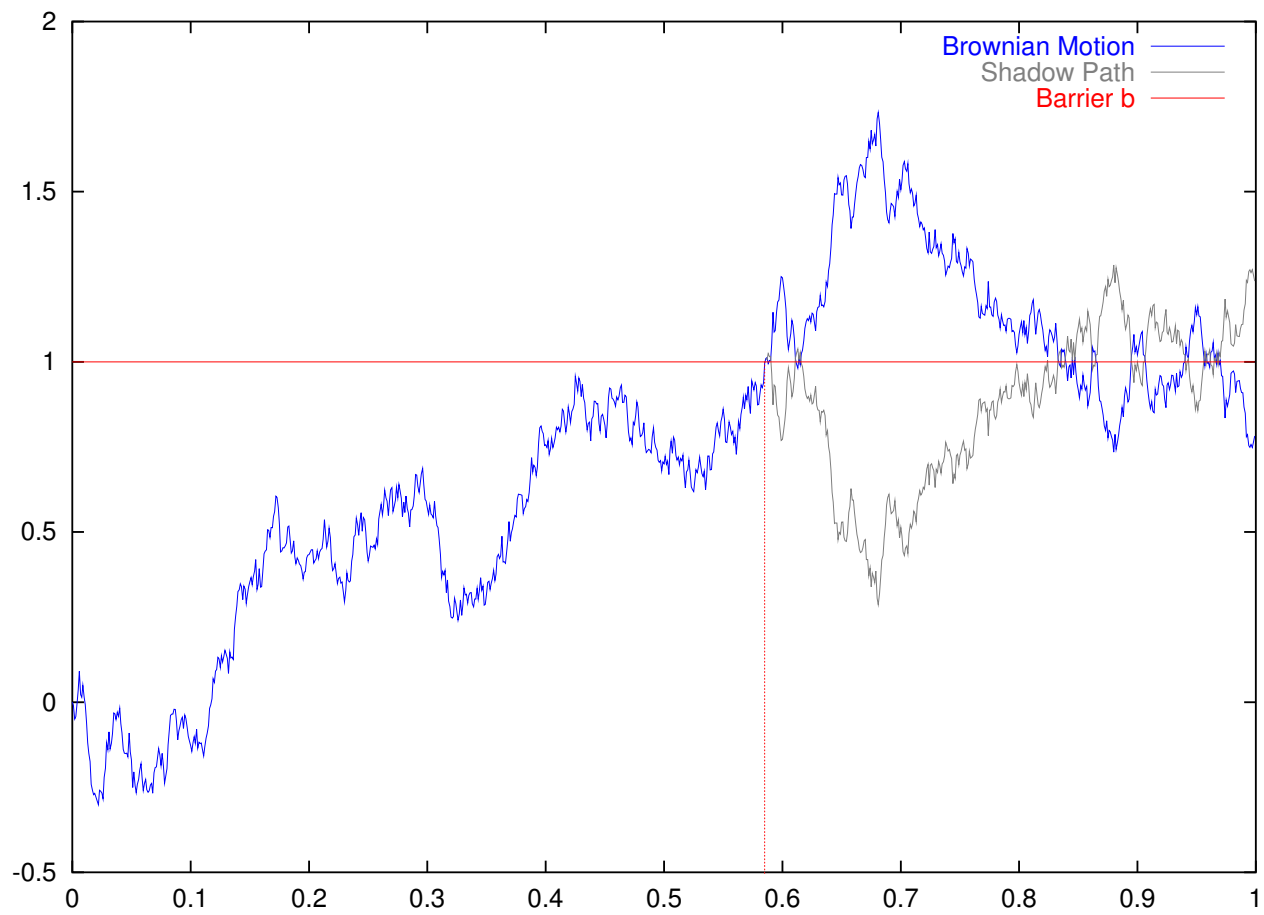
## 3.2 Reflection Principle

**Theorem 2 (Reflection Principle).** *Let  $b > 0$  and  $W_t$  a Brownian motion. Let  $T_b$  be the first passage time with respect to level  $b$  defined as*

$$T_b := \inf_{t \geq 0} (W_t = b)$$

*Then we can express the joint distribution of first passage time and Brownian motion directly with the distribution of the Brownian motion. In particular we have*

$$P(T_b < t, W(T) < x) = P(W(T) > 2b - x) \quad \forall x < b$$



A generalization of the Reflection Principle to a situation with two barriers leads to

**Lemma 3 (Bounded Green's Function).** *Let  $\rho = 0$ . Then*

$$\begin{aligned} G(\tau, X, v, X') \\ = \sum_{-\infty}^{\infty} \left( q(\tau, X' - X_n, v) - q(\tau, X' + X_n - 2a, v) \right) \end{aligned}$$

where  $X_n = X + 2n(b - a)$ .

## 3.3 Limitations

- One should note, that the theorem relies heavily on the symmetry of the Brownian Motion.
- As the lemma is a direct consequence of the theorem, we have to check if the free Green's function in our problem is indeed even in  $Y$ .
- This is only true for the case  $\rho = 0$ : Then  $q(\tau, Y, v)$  is invariant with respect to the reflections  $Y, k \rightarrow -Y, -k$ . Else the distribution is skewed and we cannot apply the lemma any longer.

## 3.4 Uncorrelated Case

In the special case with  $\rho = 0$  we can nevertheless use this method in a straightforward manner. First of all we know that  $N \equiv 0$ .

The Green's function keeps its basic form, only the functions  $\mu$  and  $\zeta$  reduce to

$$\mu(k) = -\frac{1}{2}\kappa$$
$$\zeta(k) = \frac{1}{2}\sqrt{k^2\sigma^2 + \kappa^2 + \frac{\sigma^2}{4}}$$

As seen before we can now write the solution of the pde as

$$W(\tau, v, X) = \int_{\hat{a}}^{\hat{b}} G(\tau, X, v, X') W(0, v, X') dX'$$

and after using the suitable payoff function as

$$W(\tau, v, X) = \int_{\hat{a}}^{\hat{b}} G(\tau, X, v, X') (e^{X'/2} - e^{-X'/2})^+ dX'$$

Performing the necessary substitutions yields the following result:

$$W(\tau, v, X) = \sum_{n=-\infty}^{\infty} W_n(\tau, v, X)$$

where

$$\begin{aligned} W_n(\tau, v, X) = & e^{\hat{b}/2} \mathbf{H}^+(\tau, X_n - \hat{b}, v) - \mathbf{H}^+(\tau, X_n, v) \\ & - e^{-\hat{b}/2} \mathbf{H}^-(\tau, X_n - \hat{b}, v) + \mathbf{H}^-(\tau, X_n, v) \\ & - e^{-\hat{b}/2} \mathbf{H}^+(\tau, X_n - 2\hat{a} + \hat{b}, v) + \mathbf{H}^+(\tau, X_n - 2\hat{a}, v) \\ & + e^{\hat{b}/2} \mathbf{H}^-(\tau, X_n - 2\hat{a} + \hat{b}, v) - \mathbf{H}^-(\tau, X_n - 2\hat{a}, v) \end{aligned}$$

and

$$\mathbf{H}^{\pm}(\tau, X, v) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{e^{-ikX + 2(\mathbf{A}(\tau, k) - \mathbf{B}(\tau, k)v)/\sigma^2}}{ik \pm 1/2}$$

## 3.5 Summary

- The method works only in the special case and cannot be directly generalized.
- The method yields formulas, which are computationally expensive and somehow inelegant.

## 4. Eigenfunction Expansion

## 4.1 Overview

**Second approach:** eigenfunction expansion.

**Principle:** Use separation of variables to write the solution of the partial differential equation as infinite series of its eigenfunctions. Replacing the final payoff with a Dirac  $\delta$ -function results then in the bounded Green's function.

**References:** Karlin and McGregor (1960); Lewis (2000)

## 4.2 Constant Volatility

Let us consider for the moment the classical Black-Scholes setting. At first we have to transform the basic pricing problem

$$V_t + \frac{1}{2}\sigma^2 S^2 V_{SS} + (r_f - r_d)SV_S - r_d V = 0$$

in its nondimensional form, first by substituting

$$\chi = \sigma^2 \tau$$

to get the form

$$V_\chi - \frac{1}{2}S^2 V_{SS} - \frac{(r_f - r_d)}{\sigma^2} SV_S + \frac{r_d}{\sigma^2} V = 0$$

Then perform the transformation

$$V(\chi, S, K) \rightarrow K\Theta(\chi, \zeta), \quad \zeta = S/K$$

to get the form

$$\Theta_{\chi} - \frac{1}{2}\zeta^2\Theta_{\zeta\zeta} - \frac{(r_f - r_d)}{\sigma^2}\zeta\Theta_{\zeta} + \frac{r_d}{\sigma^2}\Theta = 0$$

followed by introducing  $X = \ln(\zeta)$  to obtain

$$\Theta_{\chi} - \frac{1}{2}\Theta_{XX} + \left(\frac{1}{2} - \frac{1}{\sigma^2}(r_f - r_d)\right)\Theta_X + \frac{r_d}{\sigma^2}\Theta = 0$$

- Finally make a change of the dependent variable

$$U(\chi, X) = e^{(\frac{r}{\sigma^2} + \gamma_{\pm}^2/2)\chi + \gamma_{\pm}X} \Theta(\chi, X)$$

where

$$\gamma_{\pm} = \pm \frac{1}{2} + \frac{r_f - r_d}{\sigma^2}$$

- The transformed problem is then just the heat equation

$$U_{\chi} - \frac{1}{2}U_{XX} = 0$$

with boundary conditions

$$U(\chi, a) = 0 = U(\chi, b)$$

- We now separate of variables via

$$U(\chi, X) = e^{-\mu\chi}g(X)$$

- This results in

$$2\mu g + g_{XX} = 0$$

with adjusted boundary conditions

$$g(a) = 0$$

$$g(b) = 0$$

- After replacing  $2\mu$  by  $k^2$  we can write the well-known general solution to this simple harmonic motion equation with boundary conditions as

$$g(X) = \Phi \cos(kX) + \Psi \sin(kX)$$

- After shifting the function via  $X \rightarrow X - a$  we can use the boundary conditions to determine the coefficients:  $g(a) = 0$  implies  $\Phi = 0$  and  $g(b) = 0$  gives us a condition on  $k$  to get a nontrivial solution:

$$k = k_n = \frac{\pi n}{b - a}$$

- Now we can write the eigenvalues  $\mu_n$  simply as

$$\mu_n = \frac{1}{2}k_n^2 = \frac{\pi^2 n^2}{2(b-a)^2}$$

- Using this result we get the following form of the eigenfunctions:

$$g_n(X) = \psi_n \sin(k_n(X - a))$$

- Combining our results we get the solution of the time-dependent problem

$$U_n(\chi, X) = e^{-\mu\chi} g_n(X) = e^{-\frac{1}{2}k_n^2\chi} \psi_n \sin(k_n(X - a))$$

- Hence the general solution as linear combination of the  $U_n(\chi, X)$  has the form

$$U(\chi, X) = \sum_{n=1}^{\infty} e^{-\frac{1}{2}k_n^2\chi} \psi_n \sin(k_n(X - a))$$

**Theorem 4 (Orthogonality of the Eigenfunctions).** *Let  $\mathfrak{N} = L^2([a, b])$  be the Hilbert space of square-integrable functions on  $[a, b]$ , where the inner product is defined as*

$$\langle f, g \rangle = \int_a^b f(X)g(X)dX$$

*Then the eigenfunctions  $g_n$  are orthogonal and we have*

$$\int_a^b g_n(X)g_m(X)dX = \delta_{nm} \left( \frac{1}{2} \Psi_n^2(b-a) \right)$$

- For the remaining task of determining the coefficients we can make use of the fact that

$$u(X) = U(0, X) = \sum_{n=1}^{\infty} \psi_n \sin(k_n(X - a))$$

- In combination with the orthogonality conditions we get

$$\int_a^b g_n(X)u(X)dX = \psi_n \frac{2b-a}{2}$$

- Therefore the coefficients  $\psi_n$  can be written as

$$\psi_n = \frac{2 \int_a^b u(X) \sin(k_n(X - a))dX}{b - a}$$

- Get the Green's Function by choosing a Dirac  $\delta$ -function as payoff function  $u(X)$ :

$$\begin{aligned}\psi_n &= \frac{2 \int_a^b \delta(X - X') \sin(k_n(X - a)) dX}{b - a} \\ &= \frac{2}{b - a} \sin(k_n(X' - a))\end{aligned}$$

- Hence the Green's function has the form

$$G(\chi, X, X') = \frac{2}{b - a} \sum_{n=1}^{\infty} e^{-\frac{1}{2}k_n^2 \chi} \sin(k_n(X - a)) \sin(k_n(X' - a))$$

The solution of the pricing problem can now again be written as convolution of the Green's function with the payoff function

$$U(\chi, X) = \int_a^b G(\chi, X, X')u(X')dX'$$

## 4.3 Stochastic Volatility

- Since the same transformations are not applicable in the stochastic volatility case, we would need to find a suitable transformation if we want to proceed in analogy.
- To avoid unnecessary calculations we will therefore look directly at the free Green's function obtained in the previous chapter. We can rewrite it as

$$G(\tau, X, v, X') = \frac{1}{2\pi} \int_{-\infty}^{\infty} \left( \Psi \sin(k(X - X')) + \Theta \cos(k(X - X')) \right) \times e^{2(\mathbf{A}(\tau, k) - \mathbf{B}(\tau, k)v)/\sigma^2} dk$$

Now we can use the same arguments as previously:

- The boundary condition  $G(\tau, X, v, \hat{a}) = 0$  yields  $\Theta = 0$ .
- The second boundary condition  $G(\tau, X, v, \hat{b}) = 0$  yields

$$k = k_n = \frac{\pi n}{\hat{b} - \hat{a}}$$

- Using the orthogonality properties to obtain the coefficients  $\Psi_n$  results in  $\Psi_n = \frac{2}{\hat{b} - \hat{a}} \int_{\hat{a}}^{\hat{b}} u(X) \sin(k_n(X - \hat{a})) dX$

We then obtain the bounded Green's function by setting  $u(X) = \delta(X - X')$ :

$$G(\tau, X, v, X') = \frac{2}{\bar{b} - \hat{a}} \sum_{n=1}^{\infty} e^{2(\mathbf{A}(\tau, k) - \mathbf{B}(\tau, k)v)/\sigma^2} \times \\ \sin(k_n(X - \hat{a})) \sin(k_n(X' - \hat{a}))$$

- As before we get our transformed pricing formula via

$$W^{(DBC)}(\tau, v, X) = \int_{\hat{a}}^{\hat{b}} G(\tau, X, v, X') (e^{X'/2} - e^{-X'/2})^+ dX'$$

- This can be written as

$$W^{(DBC)}(\tau, v, X) = \sum_{n=1}^{\infty} e^{2(\mathbf{A}(\tau, k) - \mathbf{B}(\tau, k)v)\sigma^2} \varphi_n^{(DBC)} \sin(k_n \ln(F/\hat{A})) dX'$$

where

$$\varphi_n^{(DBC)} = \frac{2 \left( (-1)^{n+1} k_n (\sqrt{\hat{B}/K} - \sqrt{K/\hat{B}}) + \sin(k_n \ln(\hat{A}/K)) \right)}{(k_n^2 + 1/4) \ln(\hat{B}/\hat{A})}$$

We are now only left with the task of undoing the previous transformations. Let us therefore apply the back-transformation

$$U^{(DBC)}(\tau, S, v) = e^{-r_d\tau + X/2} KW(\tau, X, v), \quad X = \ln(F/K)$$

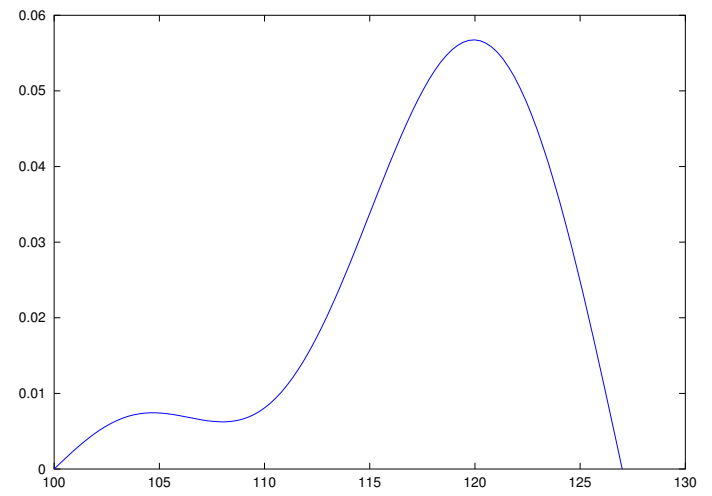
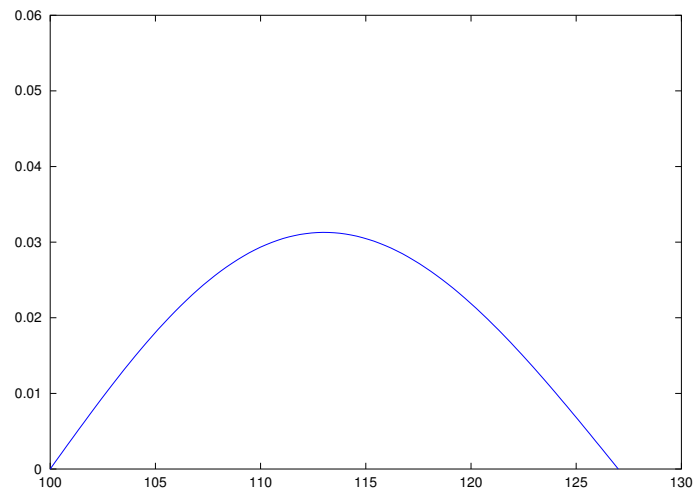
and obtain as final pricing formula

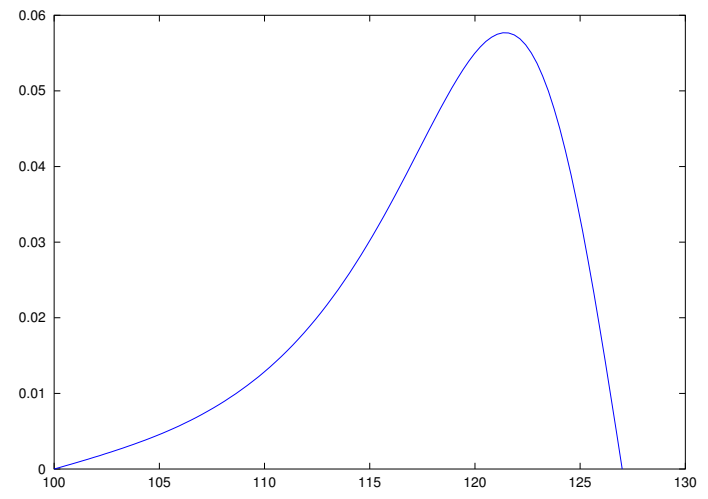
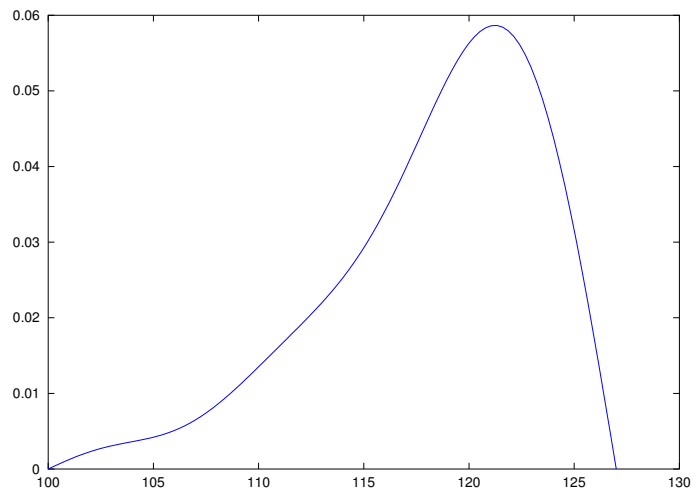
$$U^{(DBC)}(\tau, S, v) = e^{-r_d\tau} \sqrt{FK} \sum_{n=1}^{\infty} e^{2(\mathbf{A}(\tau, k_n) - \mathbf{B}(\tau, k_n)v)/\sigma^2} \times \\ \varphi_n^{(DBC)} \sin(k_n \ln(F/\hat{A}))$$

## 4.4 Limitations

Now we can see, why a generalization of Lipton's special case cannot lead to equations of the same form:

- If we introduce a correlation  $\rho \neq 0$  we would obtain a  $W_{vX}$  term in the partial differential equation. (cosine-terms!)
- Even a choice of  $r_d \neq r_f$  would cause problems, since this results in time-dependent barriers. (Separation of Variables makes no sense!)





## 4.5 Summary

- The method works only in the special case and cannot be directly generalized.
- The method yields elegant formulas, which can be easily implemented.

## 5. Application

## 5.1 Overview

We can use several methods to adopt this approach to options with a different payoff scheme:

- Solving the Integral: DNT, Digitals
- Replication Portfolios: DOT, Knock-Ins
- Asymptotic Approximation: Vanillas, Single Barriers, Power Options

## 6. Implementation

## 6.1 Overview

- The Eigenfunction Expansion Method was implemented in Mathematica 3.0
- As benchmarks we used the analytic solutions for vanillas and the finite-difference method for barriers.
- Lower and upper artificial barriers at 1 and 500 were introduced, if the need arose.

## 6.2 Parameters

Parameter	Value
$\kappa$	1.98937
$r_d$	0.036814
$r_f$	0.036814
$\sigma$	0.33147
$S_0$	123.4
$K$	120
$v_0$	0.014328
$\tau$	0.50137
$\theta$	0.011876
$\rho$	0

## 6.3 Vanillas

Type	Vanilla	Vanilla	Digital	Power
P/C	Call	Put		Call
$V_{Ana}$	5.56965	2.23072		
$V_{Eigen}$	5.56755	2.22973	0.639612	33.2701
Deviation (in %)	0.04	0.04		
Number of Iterations	302	342	164	302
Elapsed Time	0.40	0.47	0.22	1.172

## 6.4 Barriers

Type	Up and Out	Double Barrier	DNT
P/C	Call	Call	
Lower Barrier		120	120
Upper Barrier	127	127	127
$V_{FD}$	0.327446	n.a.	n.a.
$V_{Eigen}$	0.325321	0.109482	0.0317398
Number of Iterations	168	4	4
Elapsed Time	0.23	0.02	< 0.001

## 6.5 Efficiency

The method performs well, if

- the barriers are narrow
- the maturity is long (!)